## Electronic Spectrum and **d**-Orbital Energies of the Planar CuCl<sup>2-</sup> Ion **in Bis(creatininium) Tetrachlorocuprate**

By **MICHAEL A. HITCHMAN** 

*(Chemistry Department, University of Tasmania, Box 252C, Hobart, Tasmania* **7001,** *Austvalia)* 

*Summary* The low temperature, polarized electronic spectrum of the (100) crystal face of (creatininium)<sub>2</sub>CuCl<sub>1</sub> is reported and interpreted in terms of a d-orbital energy s reported and interpreted in terms of a *u*-orbital energy<br>sequence  $d_{x2-y2}$   $>$   $d_{xy}$   $>$   $d_{xz}$   $>$   $d_{yz}$   $>$   $d_{z2}$  for the planar  $CuCl<sub>4</sub><sup>2-</sup>$  ion present in this compound; the marked temperature dependence of the spectrum and its comparison with that of the similar compound (PhCH<sub>2</sub>CH<sub>2</sub>-NH,Me),CuCl, suggests that coupling both with lattice modes and an out-of-plane bonding mode of very low energy is important in the mechanism by which the electronic transitions gain intensity in compounds of this kind.

THE electronic spectrum of planar CuCl<sub>4</sub><sup>2-</sup> is of particular interest as the energy levels of this ion have been the subject of numerous theoretical studies, both of a sophisticated nature<sup>1,2</sup> and in testing simple models of the bonding in transition metal complexes.<sup>3</sup> Until recently, however, the only compound known to contain this ion with no axial ligand interaction was  $(PhCH_2CH_2NH_2Me)_2CuCl_4$ ,  $(nmph)<sub>2</sub>CuCl<sub>4</sub>$ .<sup>4</sup> The electronic spectrum of this complex suggested d-orbital energies in agreement with simple theory except that the  $d_{z2}$  orbital was *ca.* 5000 cm<sup>-1</sup> lower in energy than expected, this being rationalized in terms of the configuration interaction with the metal 4s orbital expected in a planar, truly 4-co-ordinate complex.<sup>5</sup> The spectrum was also unusual in showing extensive vibrational fine structure at low temperatures,<sup> $6,7$ </sup> some of which was attributed to coupling with what are formally lattice modes.' However, the spectral analysis was hampered by the unfortunate crystal packing of the  $CuCl<sub>4</sub><sup>2-</sup>$  groups, which did not allow the clear resolution of the crystal spectra into their molecular components. Recently, the



crystal structure of a second compound containing planar  $CuCl<sub>4</sub><sup>2-</sup> ions, (creation inium)<sub>2</sub>CuCl<sub>4</sub>, has been reported.<sup>8</sup> The$ dimensions of the CuCl $_{4}^{2-}$  unit in this compound are virtually identical to those in  $(nmph)<sub>2</sub>CuCl<sub>4</sub>$ . Moreover, ligand interactions orthogonal to the plane of the complex in  $(creation inium)<sub>2</sub>CuCl<sub>4</sub>$  are confined to two very distant oxygen atoms (Cu-0 distances **3-64** A) and the crystal packing allows the molecular spectra to be resolved quite readily. This compound therefore provides an ideal means **of** investigating the energy levels in simple systems such as  $CuCl<sub>a</sub><sup>2</sup>$ , and how these are affected by the environment of the surrounding lattice and the presence of very distant axial ligands.

Green crystals of (creatininium)<sub>2</sub>CuCl<sub>4</sub> were prepared by the method of Udopa and Krebs<sup>8</sup> and had a satisfactory analysis. The compound crystallized with the **(100)** face well developed and electronic spectra were recorded over a temperature range using a Cary **17** spectrophotometer with the electric vector of light parallel to the *b* and c crystal axes. Typical spectra are shown in the Figure.



**FIGURE.** Electronic spectrum of the (100) crystal face of (creatininium),CuCl, measured at **290** and **8** K with the electric vector of light parallel to the  $b$   $(----)$  and  $c$   $(----)$  crystal axes.

The complex crystallizes in the spacegroup *P2,/c* with the unit cell parameters  $a = 8.080$ ,  $b = 7.831$ ,  $c =$ 13.922 Å, and  $\beta = 113.77^{\circ.8}$  The CuCl<sub>4</sub><sup>2-</sup> ion in (creati $ninium)_2$ CuCl<sub>4</sub> is rigorously centrosymmetric with Cu-Cl bond lengths of **2-233** and **2.268** *8,* and Cl-Cu-C1 angles of 89.9 and  $90.1^\circ$ . The complex thus belongs to the  $D_{2h}$ point group, with the departure from  $D_{4h}$  symmetry being due to the slight difference between the Cu-C1 bond distances. **A** molecular co-ordinate system was defined with  $x$  and  $y$  parallel to the shorter and longer Cu-Cl bonds and *z* orthogonal to these. The squares of the molecular projections made by a unit vector along the *b*  and c crystal axes are  $0.98x^2 + 0.01y^2 + 0.01z^2$  and  $0.00x^2 +$  $0.04y^2 + 0.96z^2$ , respectively, so that the *b* and *c* crystal spectra (Figure) correspond almost exactly to the *x* and *z*  molecular spectra. The vibronic selection rules for a copper(II) complex with  $D_{2h}$  symmetry and a  ${}^2A_g(x^2 - y^2)$ ground state are that transitions to the  ${}^{2}B_{1g}(xy)$ ,  ${}^{2}B_{2g}(xz)$ ,

and  ${}^{2}B_{3g}(yz)$  states are forbidden in *z*, *y*, and *x* polarization, respectively (the d-orbital containing the unpaired electron being indicated in parentheses). The spectra at 8 K show peak maxima at 12,500, 13,840, and 16,530cm-l in *x*  polarization and 14,185 and 16,390 cm-l in *z* polarization. The selection rules therefore clearly suggest that the lowest energy peak is due to the transition  ${}^2B_{1g}(xy) \leftarrow {}^2A_g(x^2 - y^2)$ . Since the departure from  $D_{4h}$  symmetry is small the  ${}^{2}B_{2g}(xz)$  and  ${}^{2}B_{3g}(yz)$  states are expected to occur close together, with the latter slightly higher in energy because of the stronger ligand perturbation along the *x* axis. As the central peak occurs at  $13,840 \text{ cm}^{-1}$  in *x* polarization when the transition  ${}^2B_{2g}(xz) \leftarrow {}^2A_g(x^2 - y^2)$  is allowed, and 14,185 cm-l in *z* polarization when both this transition and that to the  ${}^{2}B_{3g}(yz)$  state are allowed, this is consistent with the  ${}^{2}B_{20}(xz)$  state lying at 13,840 cm<sup>-1</sup> and the  ${}^{2}B_{3q}(yz)$  state being at *ca*. **14**,530 cm<sup>-1</sup>. The highest energy peak is then assigned to the transition  ${}^2A_g(z^2) \leftarrow {}^2A_g(x^2 - y^2)$ .

The excited state energies of (creatininium)<sub>2</sub>CuCl<sub>4</sub> are quite similar to those in  $(nmph)<sub>2</sub>CuCl<sub>4</sub>$   $[^{2}B_{1g}(xy) = 12,500$ cm<sup>-1</sup>,  ${}^{2}B_{2g}(xz) = 14{,}050 \text{ cm}^{-1}$ ,  ${}^{2}B_{3g}(yz) = 14{,}450 \text{ cm}^{-1}$ ,  $^{2}A_{g}(z^{2})$  = 17,000 cm<sup>-1</sup>],<sup>6</sup> the observed sequence of *d*orbitals agreeing with some theoretical calculations1 but not with others.<sup>2</sup> The values  $E(xy)$  **12,500** cm<sup>-1</sup> and  $E(xz,yz)$ *ca.* **14,200** cm<sup>-1</sup> suggest angular overlap parameters  $e_{\sigma}$ *ca.* 5300 cm<sup>-1</sup> and  $e_{\pi}$  *ca.* 850 cm<sup>-1</sup> for Cl<sup>-1</sup> towards Cu<sup>II</sup> in these complexes.<sup>9</sup> The high energy of the  ${}^2A_g(z^2)$  state in planar complexes of this kind has been rationalized in terms of configuration interaction between the  $a_{1g}(z^2)$  and  $a_{1a}(4s)$  orbitals.<sup>5</sup> The magnitude of this depends on the difference in ligand interaction along the *z* and *(xy)* axes, which, because of the diffuse nature of the 4s orbital, is only expected to be significant when the axial bonds are very long or entirely absent as in  $(nmph)<sub>2</sub>CuCl<sub>4</sub>$ . In  $(\text{creatininium})$ <sub>2</sub>CuCl<sub>4</sub> the <sup>2</sup> $A_a(z^2)$  state is 600 cm<sup>-1</sup> lower in energy than that in  $(nmph)<sub>2</sub>CuCl<sub>4</sub>$ , which may indicate a weak interaction with the axial oxygen atoms (Cu-0 distance **3.64** A) in the former compound. This would correspond to a decrease of *ca. SY,* in the magnitude of the coefficient of the  $4s$  orbital in the  $a_g(z^2)$  wavefunction<sup>5</sup> of (creatininium),CuCl, compared with the value of *ca.* **0.25**  estimated<sup>7</sup> for  $(nmph)<sub>2</sub>CuCl<sub>4</sub>$ .

The low-temperature spectrum of (creatininium),CuCl, differs markedly from that of  $(nmph)<sub>2</sub>CuCl<sub>4</sub>$  in that it contains virtually no evidence of the extensive vibrational fine structure observed<sup>6,7</sup> for the latter compound. Only the lowest energy peak shows a barely discernible progression, the spacing of  $ca$ . 260 cm<sup>-1</sup> being similar to those observed in  $(nmph)<sub>2</sub>CuCl<sub>4</sub>$  and suggesting that it involves excitation of the *alg* symmetric stretching mode. Since the geometries of the CuCl<sub>4</sub><sup>2-</sup> groups in the two compounds are virtually identical, this agrees with the supposition<sup>7</sup> that in  $(nmph)<sub>2</sub>CuCl<sub>4</sub>$  coupling with what are formally lattice modes plays an important part in the vibronic intensity mechanism. As with  $(nmph)<sub>2</sub>CuCl<sub>4</sub>$ , the spectral peaks of  $(creation inium)<sub>2</sub>CuCl<sub>4</sub>$  show marked decreases in intensity on cooling, and significant shifts of the band maxima to higher energy. The  ${}^2A_a(z^2) \leftarrow {}^2A_a(x^2 - y^2)$  transition is quite exceptional in this respect, the band maximum shifting by  $ca. 1000 \text{ cm}^{-1}$ . The crystal structure indicated an anomalously large temperature factor perpendicular to the CuCl<sub>4</sub><sup>2-</sup> plane for the chlorine atoms in (creatininium)<sub>2</sub>- $CuCl<sub>4</sub>$ , suggesting either an out-of-plane bending mode of very low energy or a slight distortion towards a tetrahedral geometry, possibly dynamic in nature.8 The excited state energies suggest that no tetrahedral distortion is present, at least at 8 K. The marked temperature dependence is consistent with the out-of-plane bending made of  $b_{2u}$  symmetry in the  $D_{4h}$  point group being active in inducing intensity and having a very low energy (much of the intensity in the spectrum of  $(nmph)<sub>2</sub>CuCl<sub>4</sub>$  was also found to be derived from this mode, the energy in this case being  $ca. 100 \text{ cm}^{-1}$ ).<sup>7</sup> It is hoped to clarify this point by making a detailed study of the electronic spectra of several crystal faces of (creatininium)<sub>2</sub>CuCl<sub>4</sub> in the near future.

The financial support of the Australian Research Grants Committee is gratefully acknowledged, as is the loan of a Cryodyne 21 cryostat from the Central Science Laboratory of the University of Tasmania.

## *(Received* **11th** *July* **1979;** *Corn.* **747.)**

- P. Ros and G. C. **A.** Schuit, *Theor. Chim. Acta,* **1966, 4, 1** ; **J** Demuynek, **A.** Veillard, and **U.** Wahlgren, *J. Amer. Chem. Soc.,*  **1973, 95, 5563.**
- *ibid.,* **1971, 8, 327;** D. W. Smith, *J.C.S. Dalton,* **1973, 1853.**  <sup>2</sup> J. Demuynek and A. Veillard, *Chem. Phys. Letters*, 1970, 6, 204; N. J. Trappeniers, G. de Brouchière, and C. A. Ten Seldam,
	- For a recent discussion of chlorocuprate(I1) complexes see D. W. Smith, *Co-ordination Chem. Rev.,* **1976, 21, 93.**
	- **R. L.** Harlow, W. J. Wells, **111,** G. W. Watt, and S. H. Simonsen, *Inorg. Chem.,* **1974, 13, 2106.**  D. W. Smith, *Inorg. Chim. Ada,* **1977, 22, 107.**
	-
	- P. Cassidy and **M. A.** Hitchman, *J.C.S. Chern. Comm.,* **1975, 837.**  M. **A.** Hitchman and P. Cassidy, *Inorg. Chem.,* **1979, 18, 1745.**
	-
	- \* **M.** R. Udopa and B. Krebs, *Inorg. Chim. Ada,* **1979, 33, 241.**
	- D. W. Smith, *Structure and Bonding,* **1978, 35,** *87.*